

Resonant conversion of QCD axions into hidden axions and suppressed isocurvature perturbations



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In collaboration with **Fuminobu Takahashi**



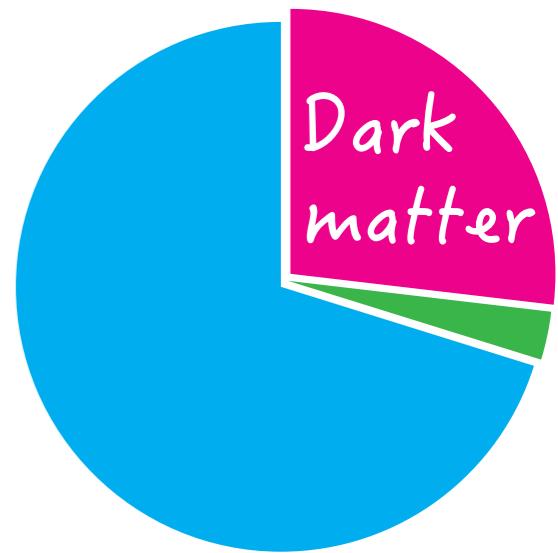
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To appear on arXiv:1411.XXXX

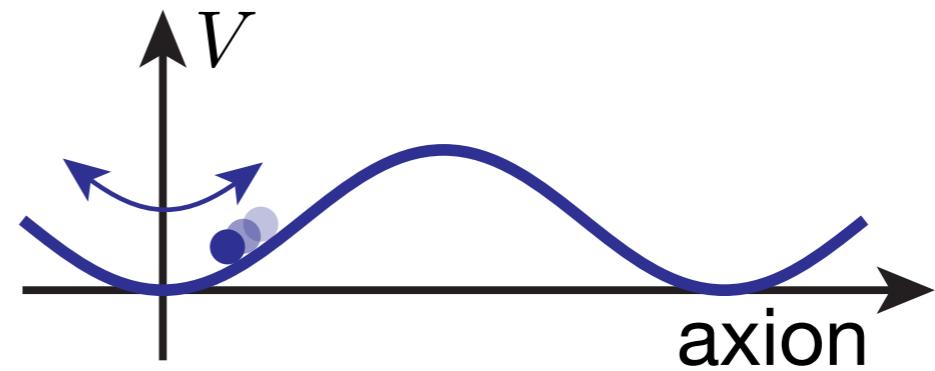
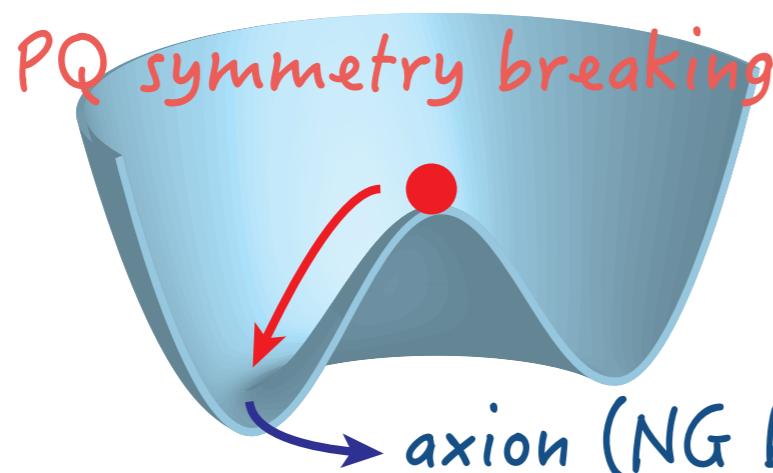
4th KIAS Workshop on Particle Physics and Cosmology
30 Oct. 2014

1. Introduction

- Our Universe is filled with the Dark matter



- Axion may exist <= strong CP problem

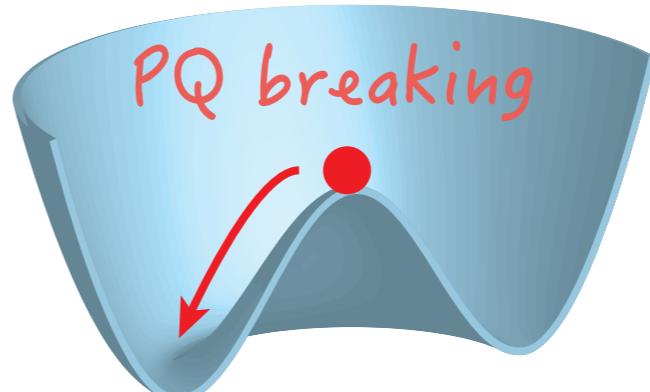
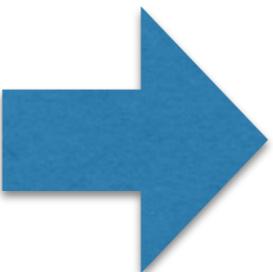


$$\Omega_a h^2 \simeq 0.195 \theta_i^2 f(\theta_i) \left(\frac{F_a}{10^{12} \text{ GeV}} \right)^{1.184}$$

Axion = Dark matter

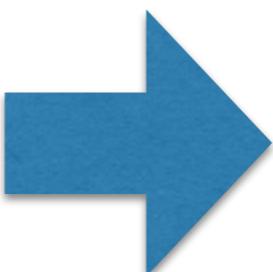
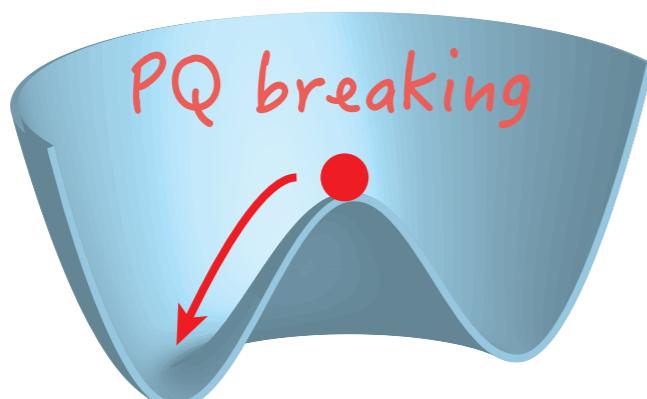
Two scenarios :

(a) inflation => PQ breaking

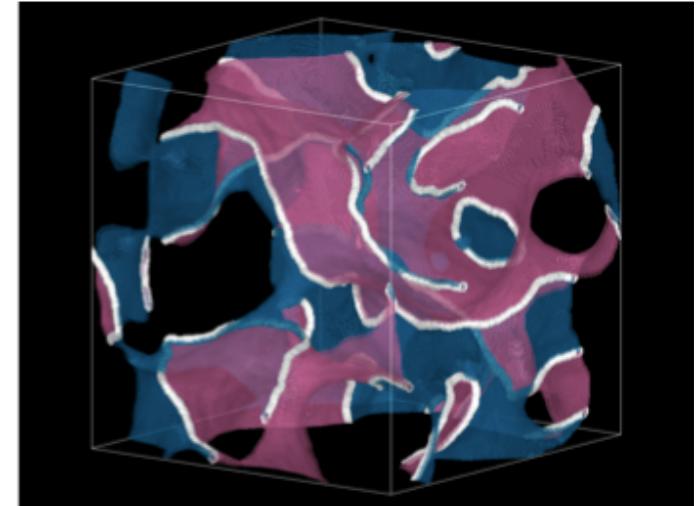


! Domain wall problem

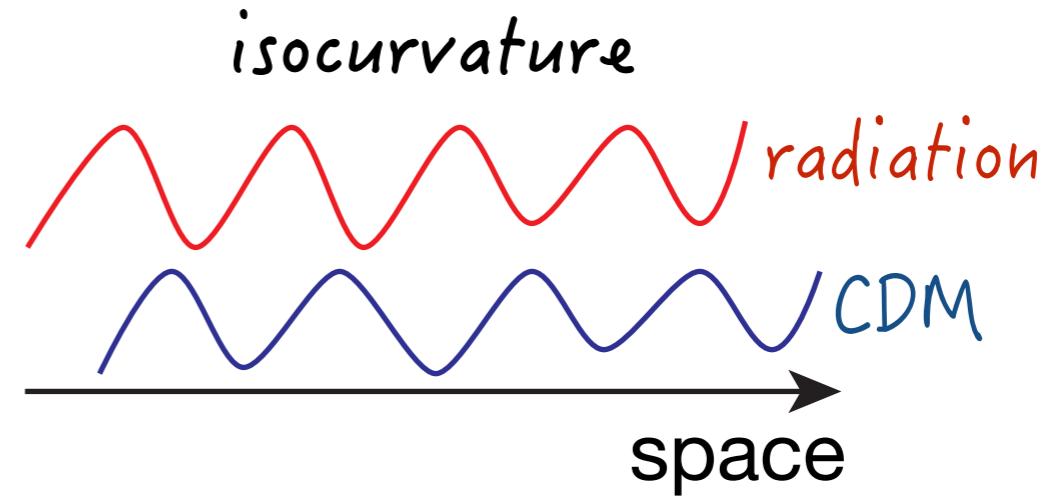
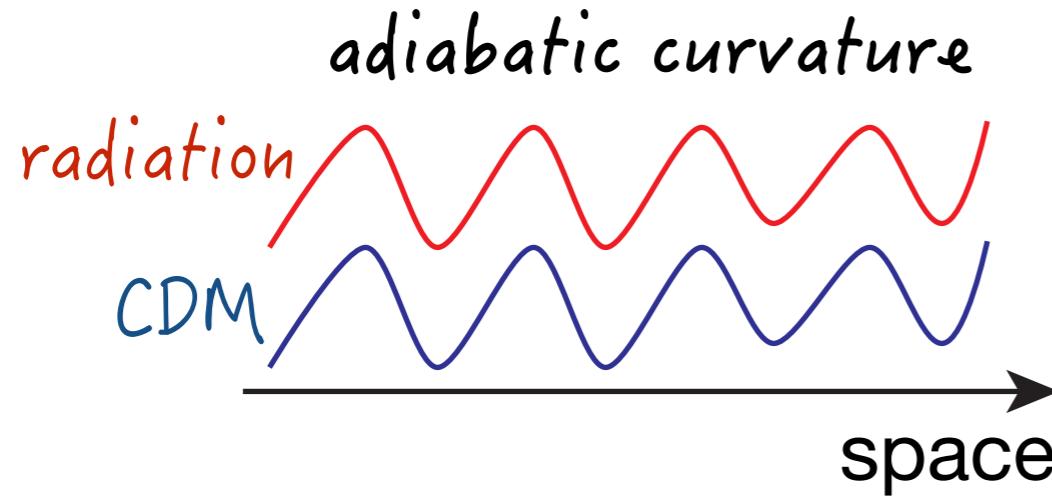
(b) PQ breaking => inflation



! Isocurvature problem



Isocurvature problem



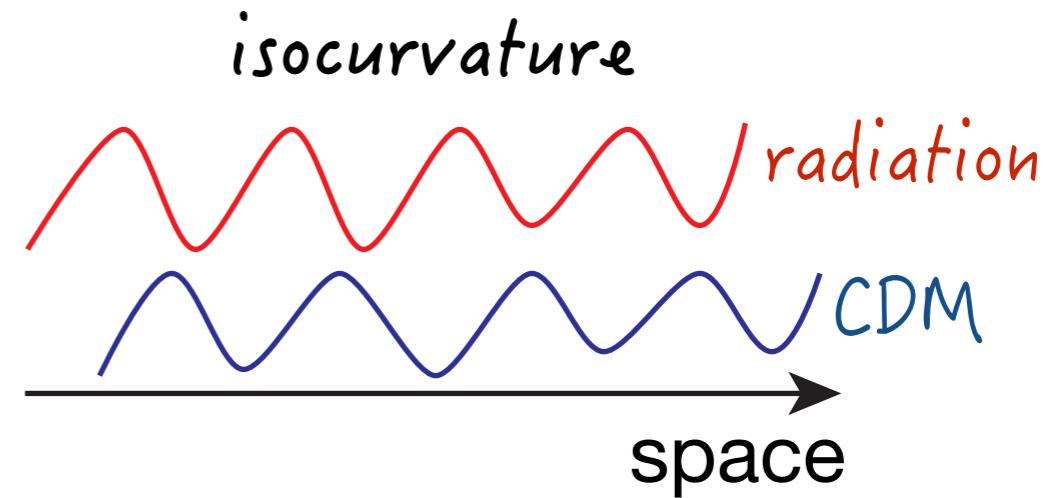
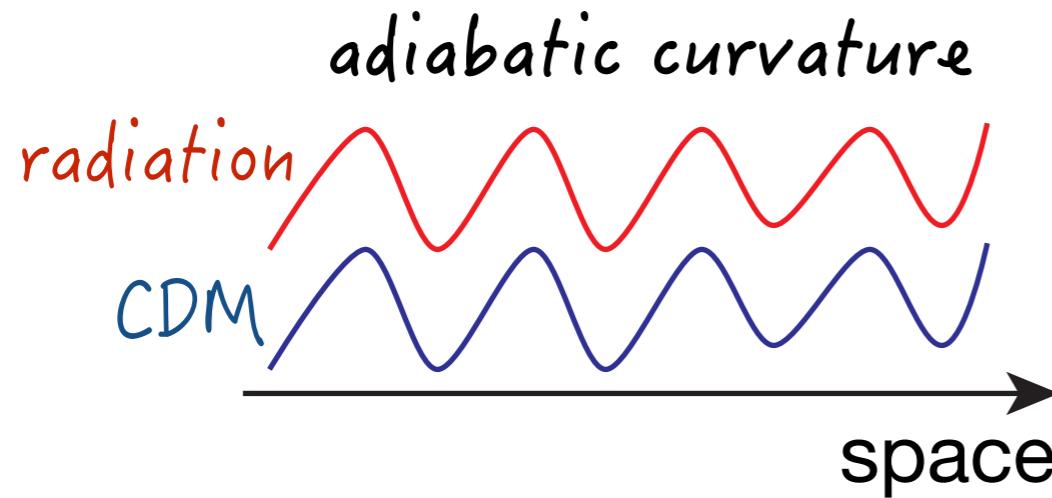
Axion CDM isocurvature perturbation

$$\Delta_{\mathcal{S}, \text{CDM}} = \left(\frac{\Omega_a}{\Omega_{\text{CDM}}} \right) \Delta_{\mathcal{S}, a} \quad \text{and} \quad \Delta_{\mathcal{S}, a} = \frac{\partial \ln \Omega_a}{\partial \theta_i} \frac{H_{\text{inf}}}{2\pi F_a} \sim \frac{\delta \Omega_a}{\Omega_a}$$

Current constraint : $\Delta_{\mathcal{S}, \text{CDM}}^2 < \frac{\beta}{1 - \beta} \Delta_{\mathcal{R}}^2$ with $\beta = 0.039$

Planck collaboration: 1303.5076, 1303.5082

Isocurvature problem



Axion CDM isocurvature perturbation

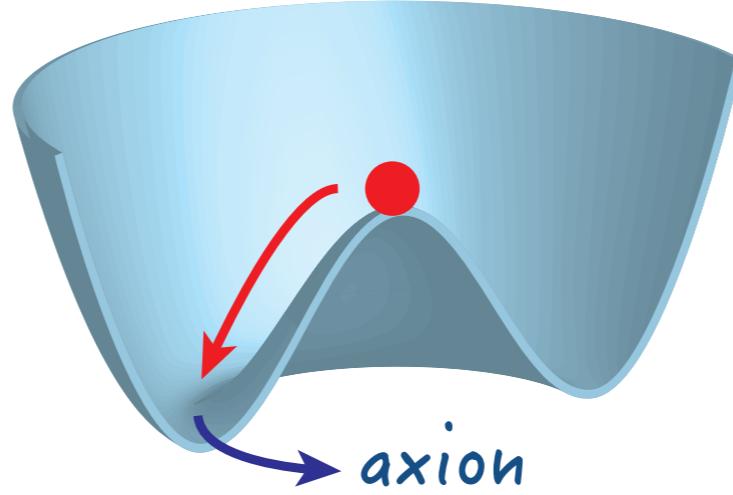
$$\Delta_{S, \text{CDM}} = \left(\frac{\Omega_a}{\Omega_{\text{CDM}}} \right) \Delta_{S,a} \quad \text{and} \quad \Delta_{S,a} = \frac{\partial \ln \Omega_a}{\partial \theta_i} \frac{H_{\text{inf}}}{2\pi F_a} \sim \frac{\delta \Omega_a}{\Omega_a}$$

Tightly constrained

The equation shows the relationship between the isocurvature perturbation in the Axion CDM model ($\Delta_{S, \text{CDM}}$) and the isocurvature perturbation in the axion field ($\Delta_{S,a}$). The ratio $\Omega_a / \Omega_{\text{CDM}}$ is highlighted with a yellow circle. The expression for $\Delta_{S,a}$ includes the partial derivative of the logarithm of the axion density with respect to a parameter θ_i , the infinite horizon value of the Hubble rate (H_{inf}), and the frequency of the mode ($2\pi F_a$). The final term $\delta \Omega_a / \Omega_a$ is also highlighted with a yellow circle.

Current constraint : $\Delta_{S, \text{CDM}}^2 < \frac{\beta}{1 - \beta} \Delta_{\mathcal{R}}^2$ with $\beta = 0.039$

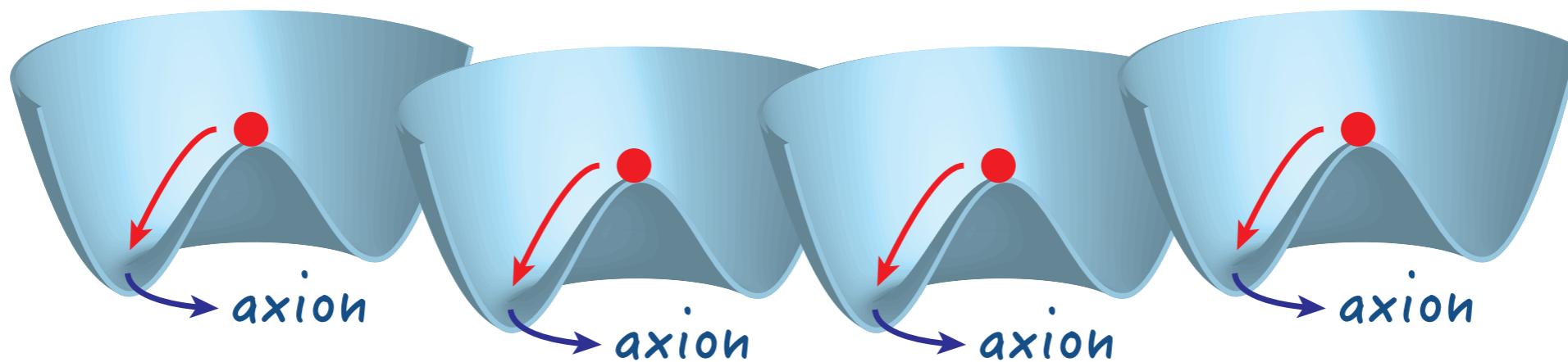
Planck collaboration: 1303.5076, 1303.5082



Single axion?

or

Multi-axion?



QCD axion + Hidden axion
with mass mixing



MSW-like resonance



Abundance is suppressed

Hill and Ross, Nucl.Phys.(1988)

QCD axion + Hidden axion
with mass mixing



MSW-like resonance



Abundance is suppressed

Hill and Ross, Nucl.Phys.(1988)

Our work

- Precise numerical calculations
- Implication for isocurvature perturbations

***New suppression mechanism
for isocurvature perturbations***

2. Model

PQ field & hidden PQ field with heavy quarks

$$\mathcal{L} = \kappa \Phi Q \bar{Q} + \frac{\lambda}{M_P} \Phi \Phi_H Q_H \bar{Q}_H$$



Potential for QCD and hidden axions

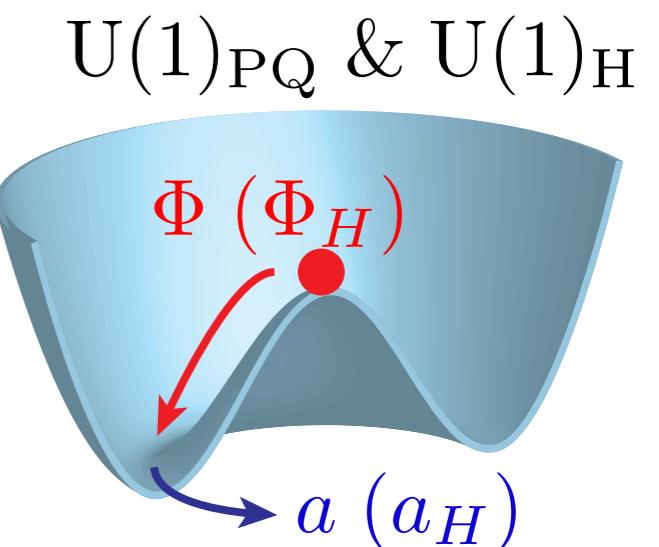
$$V(a, a_H) = m_a^2(T) F_a^2 \left[1 - \cos \left(\frac{a}{F_a} \right) \right] + m_H^2 F_H^2 \left[1 - \cos \left(\frac{a_H}{F_H} + \frac{a}{F_a} \right) \right]$$

m_a, m_H : masses

F_a, F_H : decay constants

Charge assignments

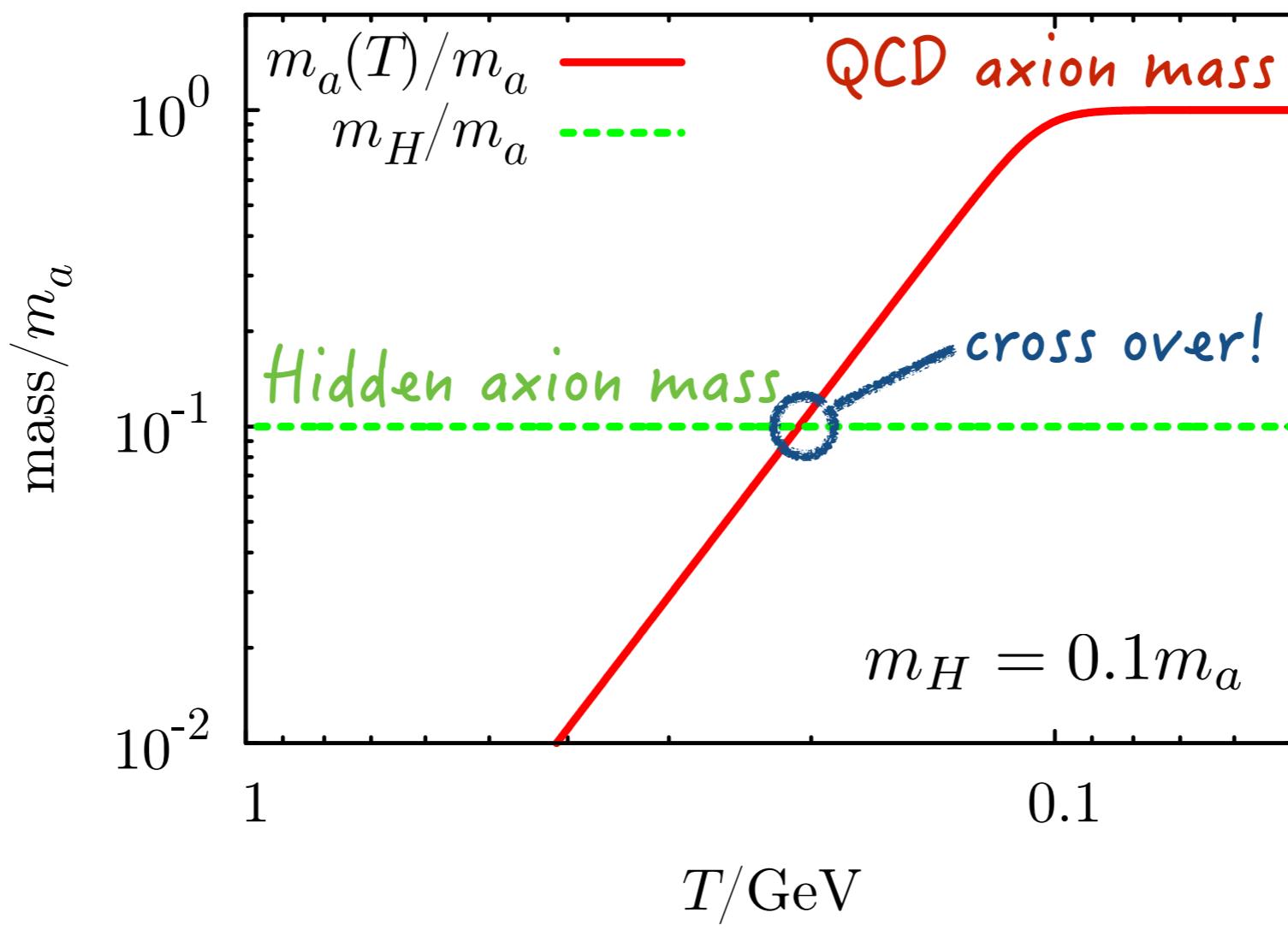
	Φ	Φ_H	Q	\bar{Q}	Q_H	\bar{Q}_H
$U(1)_{PQ}$	1	0	1/2	-1/2	1/2	-1/2
$U(1)_H$	0	1	0	0	1/2	-1/2



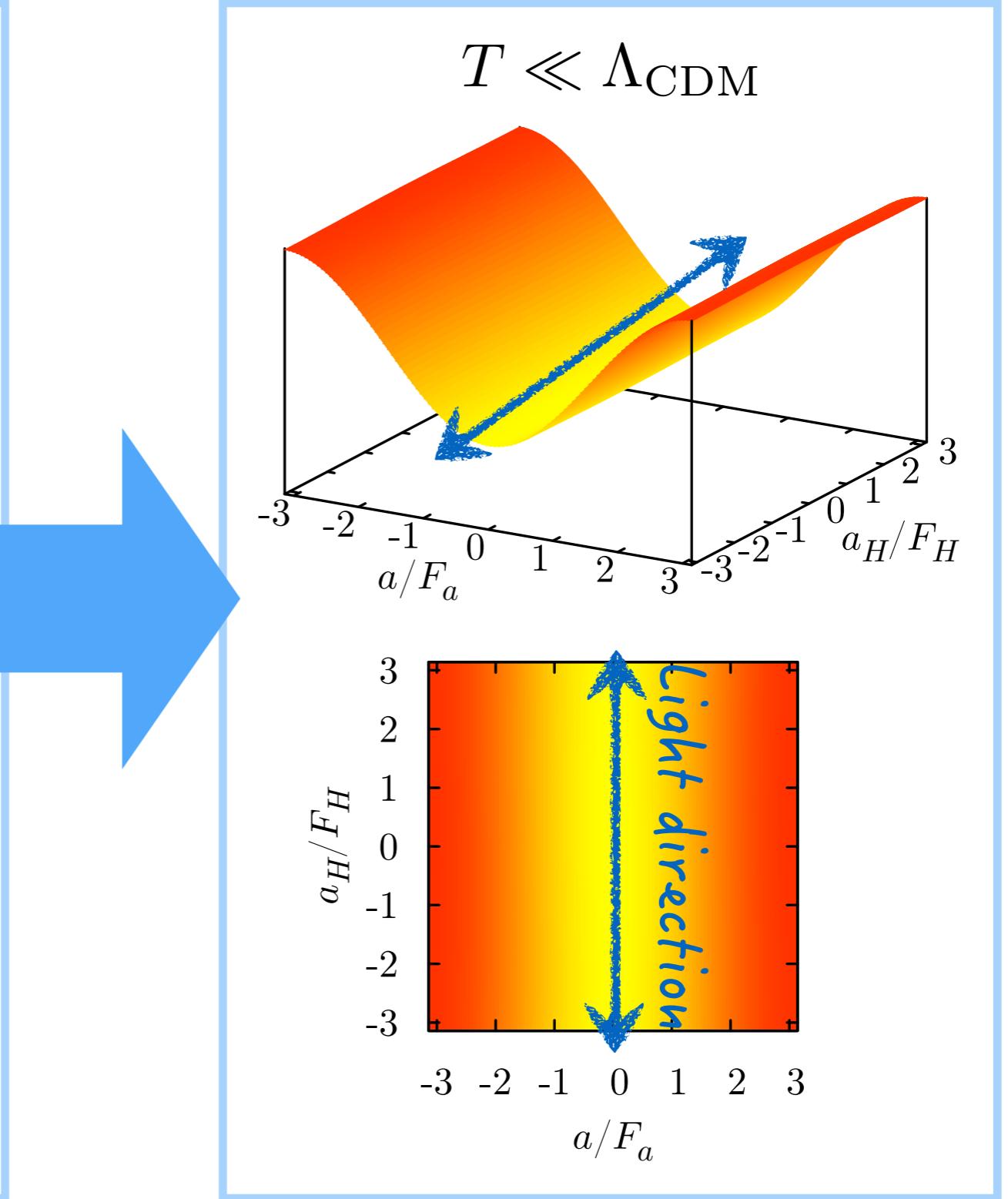
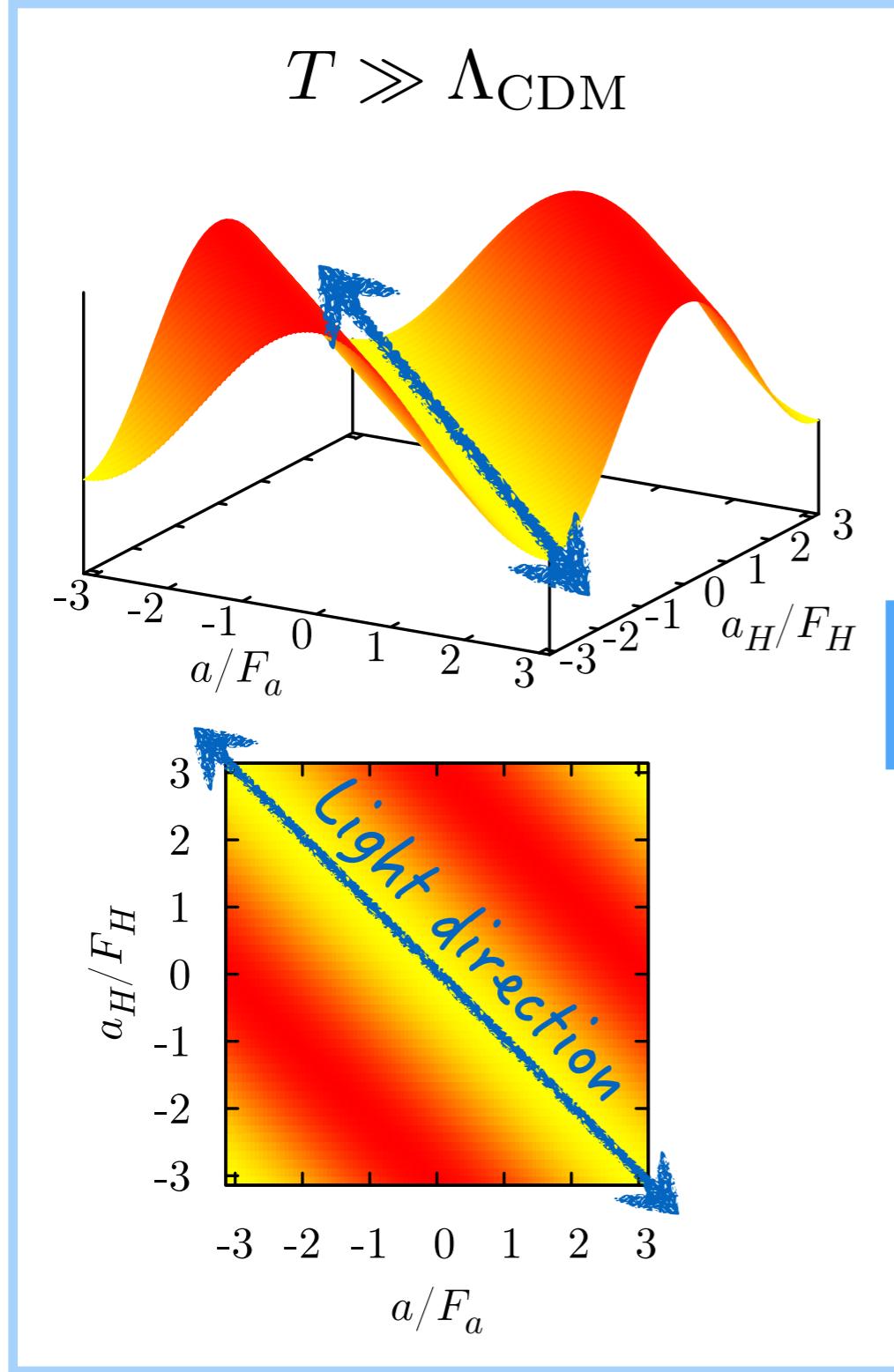
Temperature dependent QCD axion mass

$$m_a(T) = \begin{cases} 4.05 \times 10^{-4} \frac{\Lambda_{\text{QCD}}^2}{F_a} \left(\frac{T}{\Lambda_{\text{QCD}}} \right)^{-3.34} & \text{for } T > 0.26\Lambda \\ 3.82 \times 10^{-2} \frac{\Lambda_{\text{QCD}}^2}{F_a} & \text{for } T < 0.26\Lambda_{\text{QCD}}, \end{cases}$$

If zero temp. QCD axion mass >> hidden axion mass...



$$V(a, a_H) = m_a^2(T) F_a^2 \left[1 - \cos \left(\frac{a}{F_a} \right) \right] + m_H^2 F_H^2 \left[1 - \cos \left(\frac{a_H}{F_H} + \frac{a}{F_a} \right) \right]$$

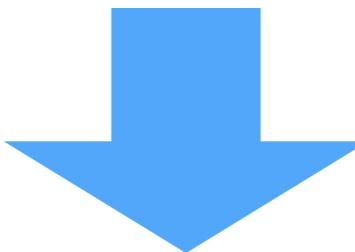


Linearized equation of motion

$$\ddot{A} + 3H\dot{A} + M^2 A = 0$$

with

$$A = \begin{pmatrix} a \\ a_H \end{pmatrix} \quad \text{and} \quad M^2 = \begin{pmatrix} m_a^2(T) + (F_a/F_H)^2 m_H^2 & (F_H/F_a)m_H^2 \\ (F_H/F_a)m_H^2 & m_H^2 \end{pmatrix}$$



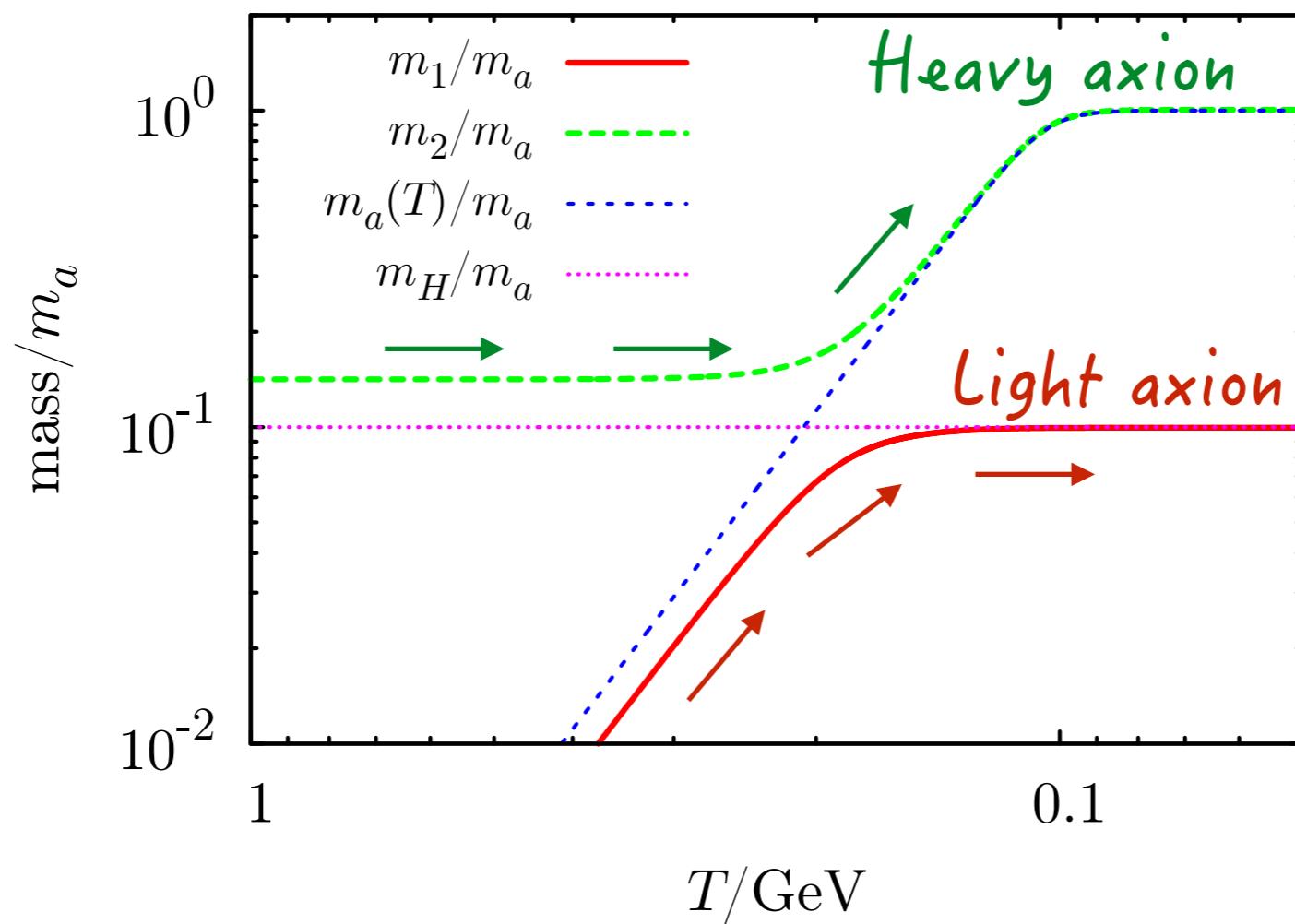
Mass eigenstates

$$\begin{pmatrix} m_1^2 & 0 \\ 0 & m_2^2 \end{pmatrix} = O^T M^2 O \quad \text{and} \quad \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = O^T A$$

a_1 : Light axion

a_2 : Heavy axion

Evolution of mass eigenvalues



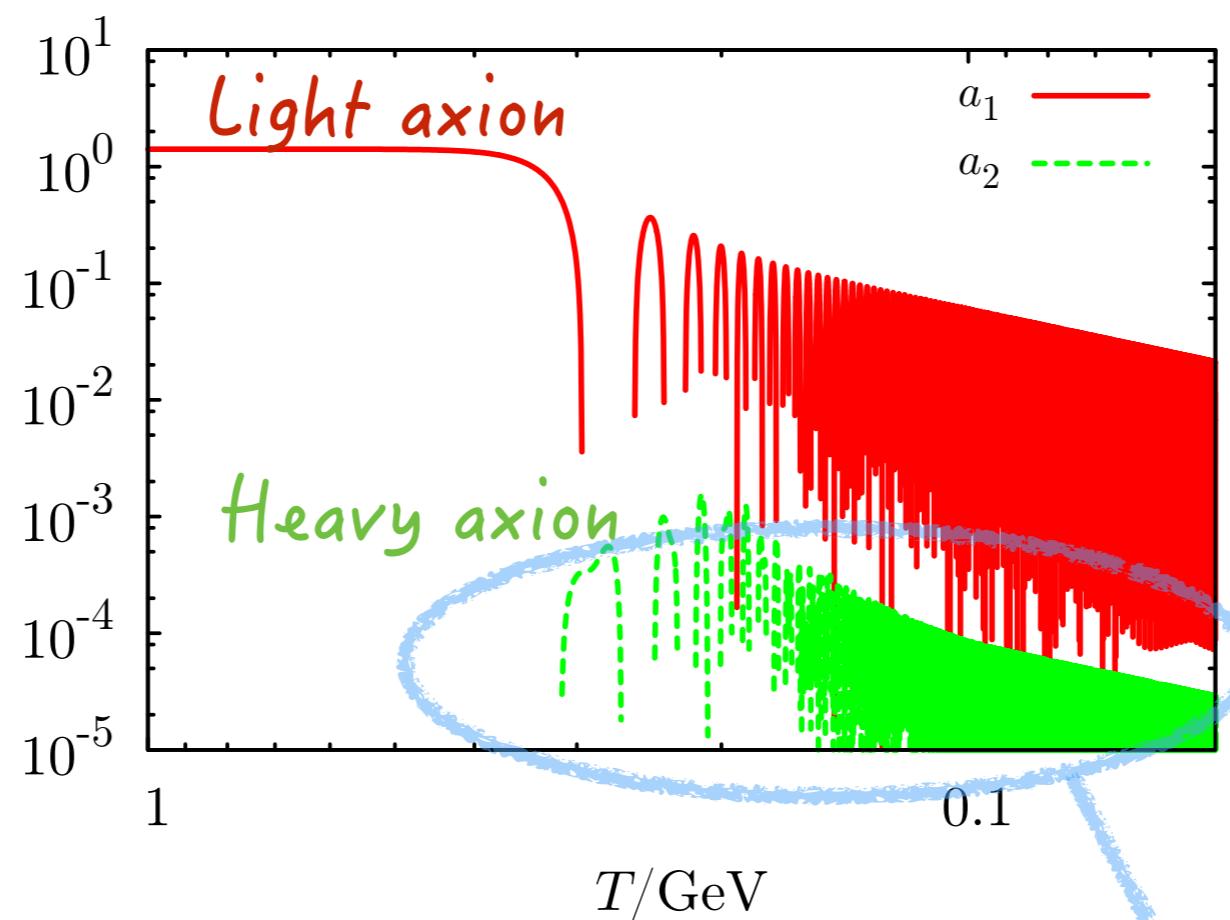
Initial condition

$$V(a, a_H) = m_a^2(T) F_a^2 \left[1 - \cos \left(\frac{a}{F_a} \right) \right] + m_H^2 F_H^2 \left[1 - \cos \left(\frac{a_H}{F_H} + \frac{a}{F_a} \right) \right] \equiv \delta = 0$$

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Heavy axion is suppressed, but...

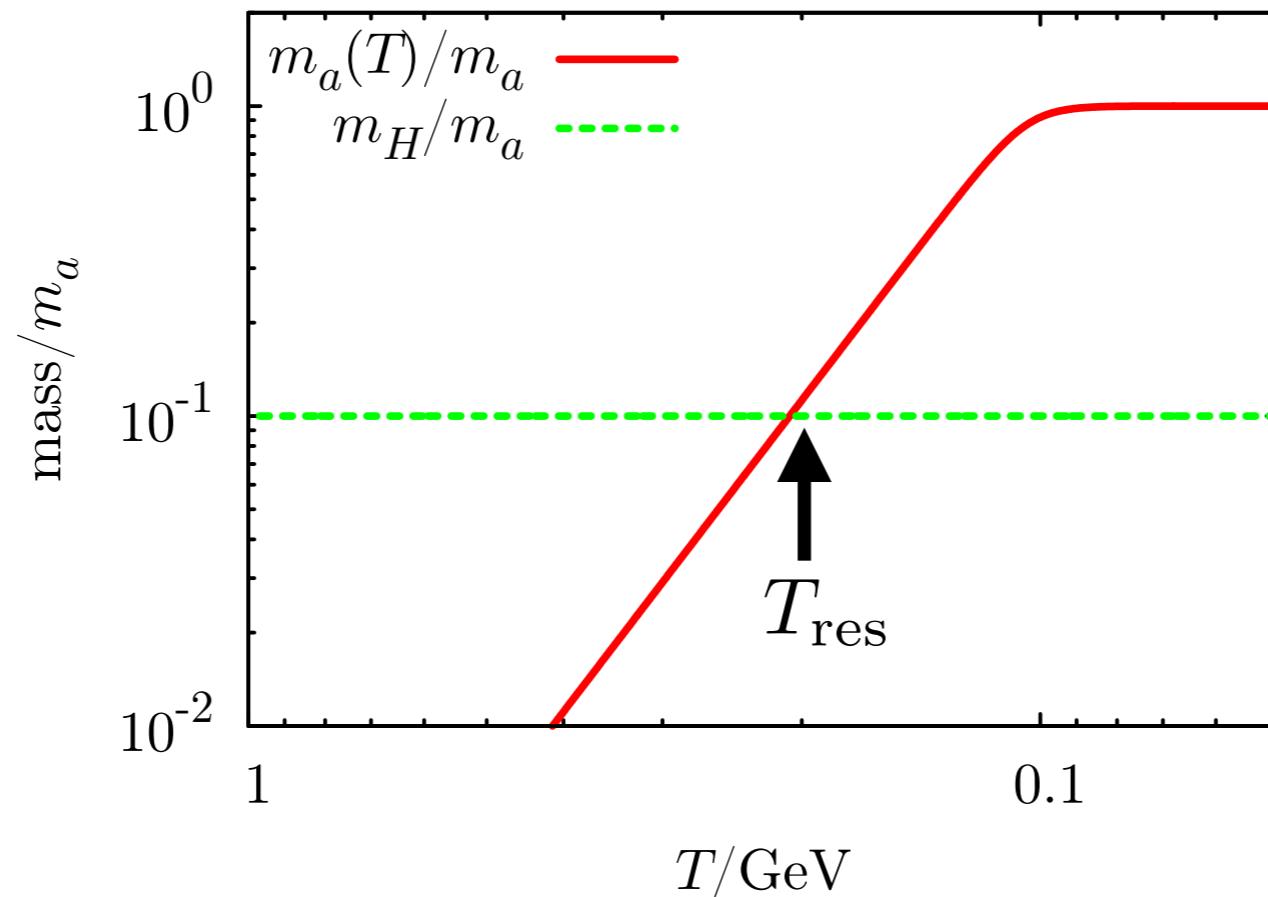


Heavy axion is suppressed

Adiabaticity, Anharmonicity

$$\xi = \frac{H(T_{\text{res}})}{m_H} \simeq 4.4 \left(\frac{m_H}{m_a} \right)^{-1.6} \frac{F_a}{M_P}$$

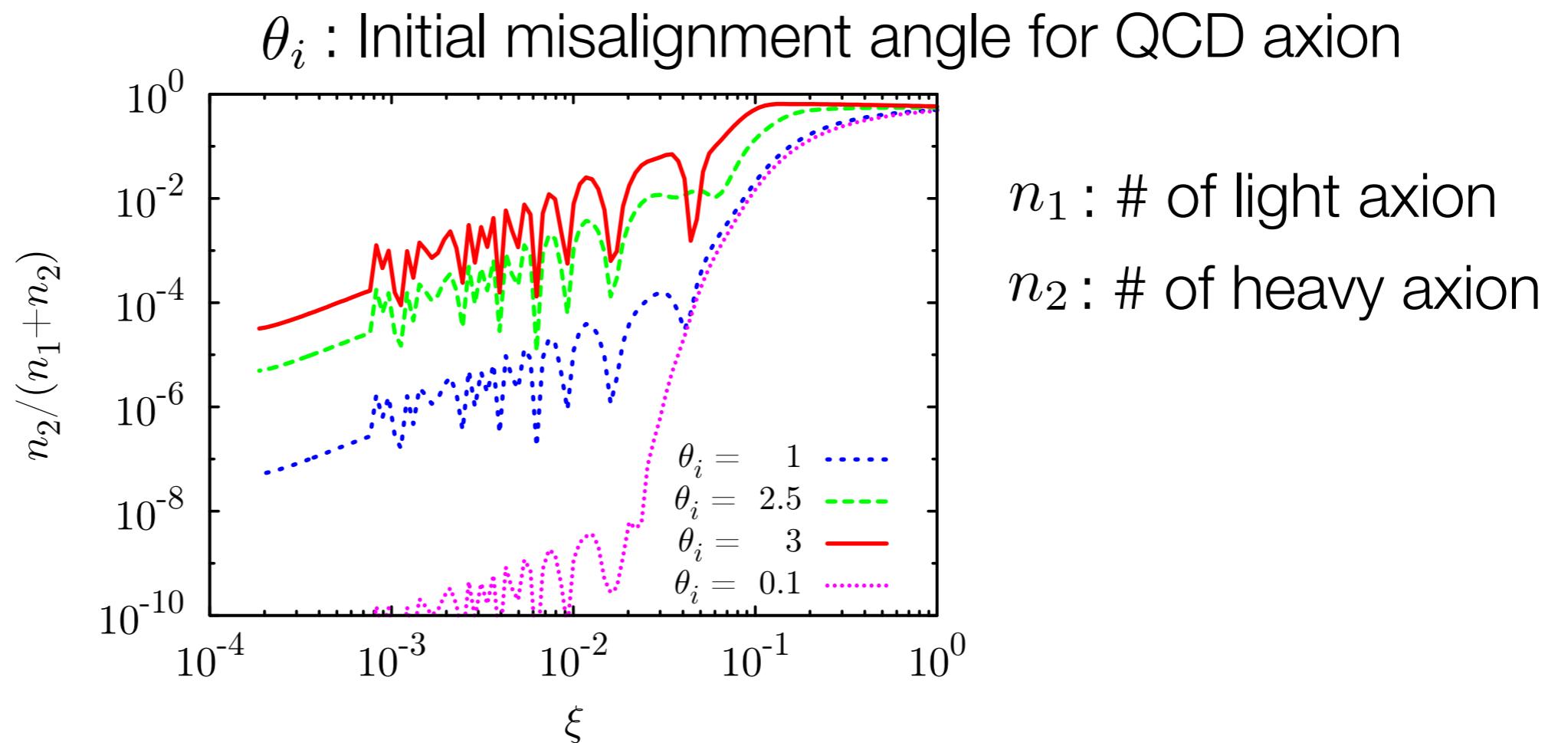
$$T_{\text{res}} \simeq 0.1 \left(\frac{\Lambda_{\text{QCD}}^2}{F_a m_H} \right)^{0.3}$$



Adiabaticity, Anharmonicity

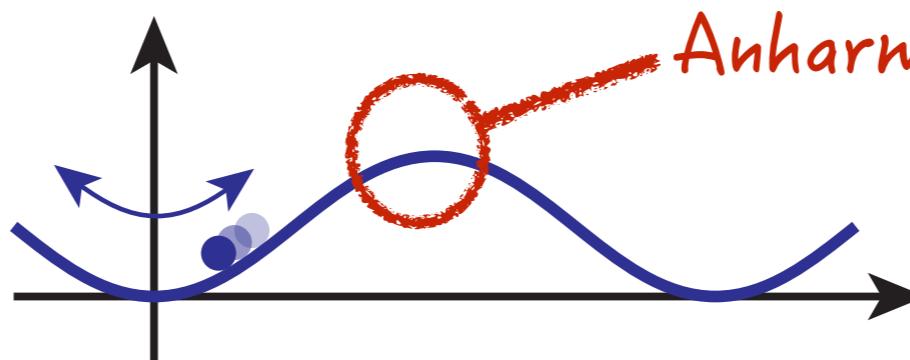
$$\xi = \frac{H(T_{\text{res}})}{m_H} \simeq 4.4 \left(\frac{m_H}{m_a} \right)^{-1.6} \frac{F_a}{M_P}$$

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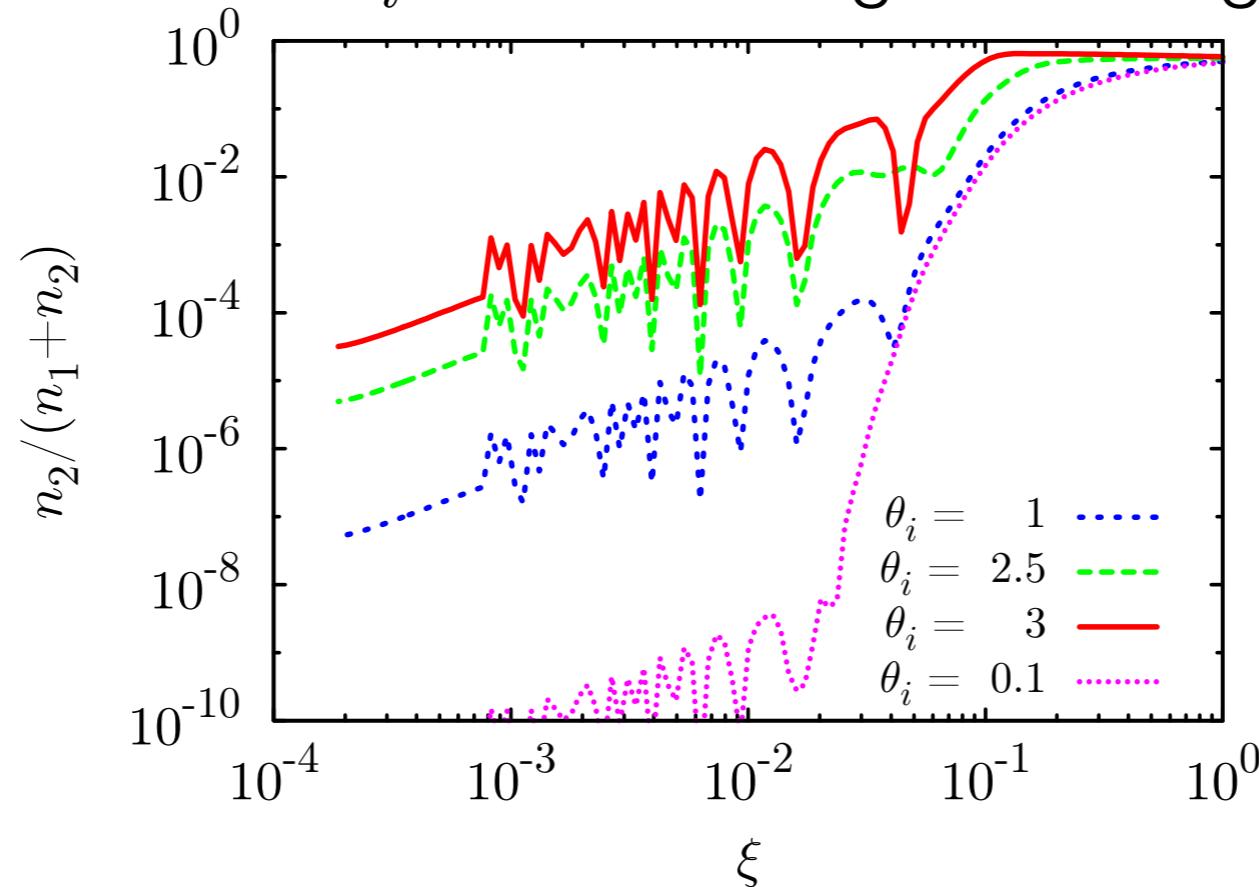


Adiabaticity, Anharmonicity

Anharmonicity = departure from the quadratic potential



θ_i : Initial misalignment angle for QCD axion

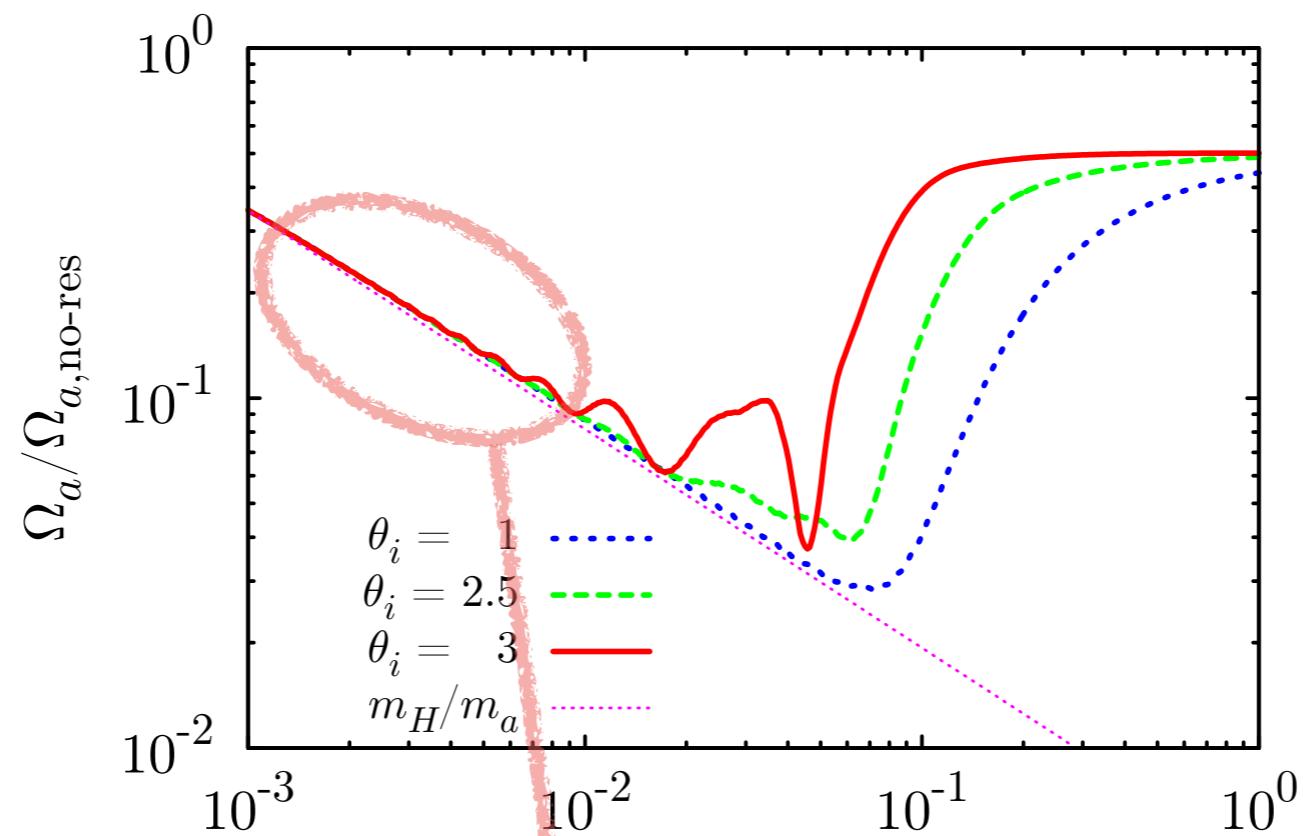


n_1 : # of light axion

n_2 : # of heavy axion

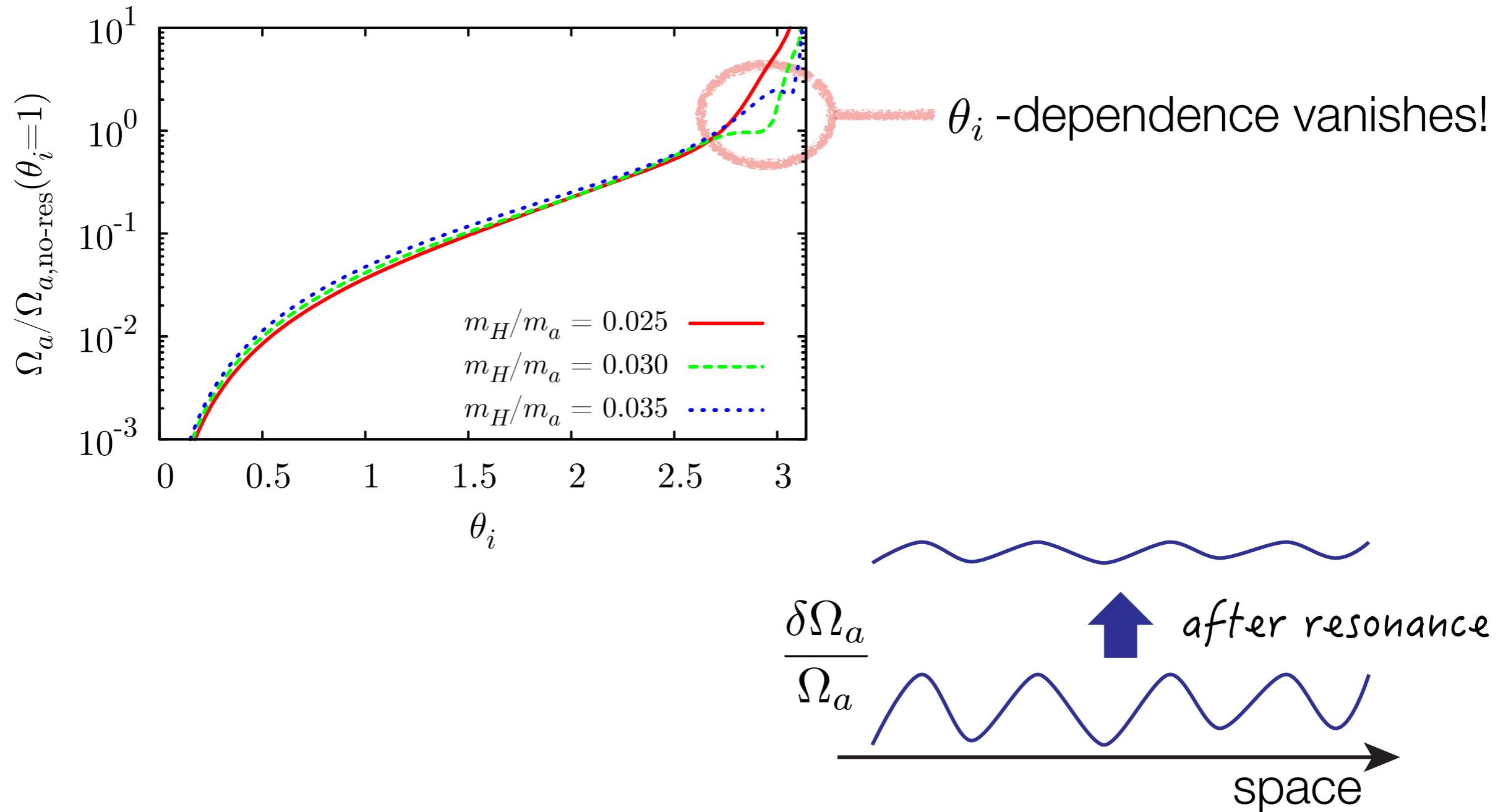
3. Cosmological implications

(i) Abundance

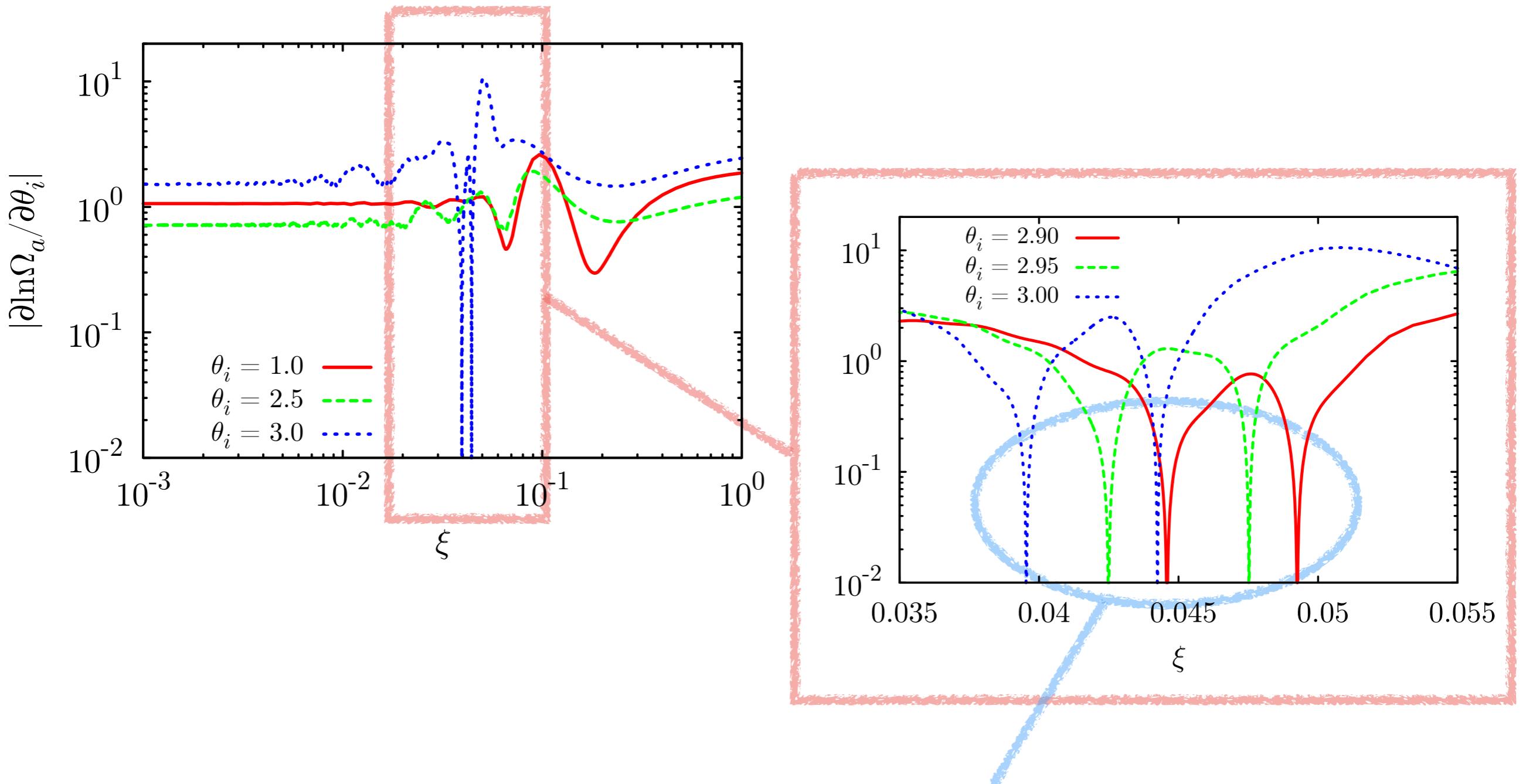


$$\Omega_a \xrightarrow{\xi} \left(\frac{m_H}{m_a} \right) \Omega_a$$

(ii) Isocurvature perturbations



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Isocurvature perturbation can be suppressed!

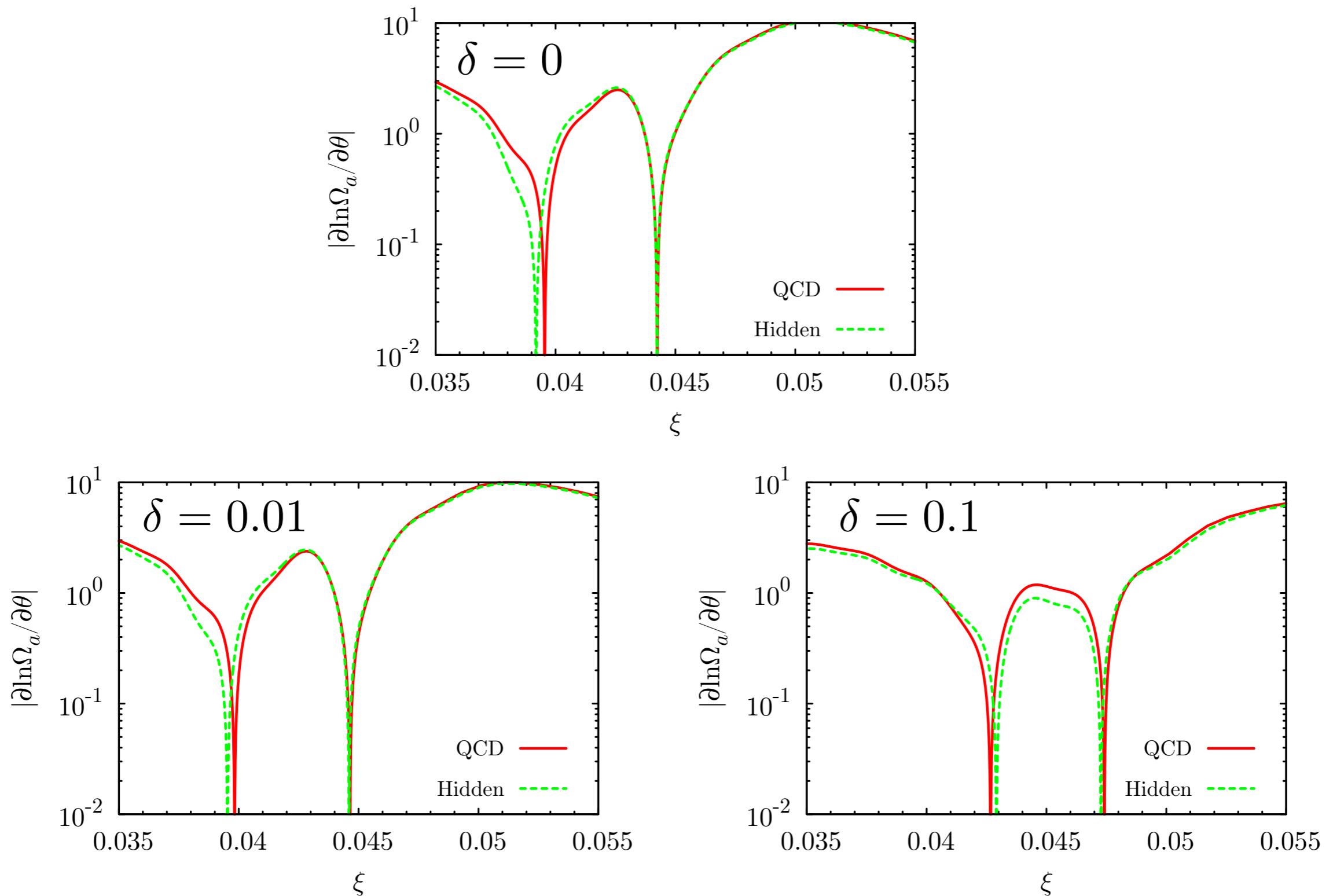
4. Conclusions

- We considered QCD axion & hidden axion with mass mixing
- MSW-like resonance takes place
and QCD axion is converted into hidden axions
- Abundance can be suppressed
- Isocurvature perturbation can also be suppressed!

**We proposed a completely new mechanism
to suppress isocurvature perturbations**

Thank you!

Isocurvature perturbation from QCD + hidden axion



Non-Gaussianity from isocurvature perturbation

$$f_{\text{NL}}^{(\text{iso})} = -1 + \frac{1}{r} \left(1 + \left(\frac{\partial^2 \ln \Omega_a}{\partial \theta_i^2} \right) \left(\frac{\partial \ln \Omega_a}{\partial \theta_a} \right)^{-2} \right)$$
$$\left(\frac{\beta}{1-\beta} \right)^2 f_{\text{NL}}^{(\text{iso})} = 40 \pm 66$$

Hikage,Kawasaki,Sekiguchi & Takahashi:1211.1095

